

Evaluation of Slow-Mode Damping Factor

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In linear aeroballistics, the kinetic energy loss due to yaw-drag ΔE_α at a coning angle-of-attack α over the period T_2 of each “slow-mode” coning cycle can be written as

$$\Delta E_\alpha(T_2) = q \cdot S \cdot V \cdot \alpha^2 \cdot C_{D_\alpha} \cdot T_2.$$

We can also formulate the instantaneous kinetic energy E_c of the orbital coning motion itself as

$$\begin{aligned} E_c &= (m/2) \cdot (r \cdot \omega_2)^2 \\ &= (m/2) \cdot (D \cdot \sin \alpha \cdot \omega_2)^2 \\ &= (m/2) \cdot (D \cdot \omega_2)^2 \cdot \alpha^2 \end{aligned}$$

where $r = D \cdot \sin \alpha \approx D \cdot \alpha$ is the coning radius of the CG of the bullet orbiting around a “mean CG” location moving smoothly along the “mean trajectory” of the bullet at its “mean velocity,” and D is the slowly varying coning distance of the CG of the bullet from its coning apex, each given in **feet**.

Now, as the coning angle α decreases (due to frictional damping for any **dynamically stable** projectile) from its initial value α_0 to its final value α_1 at the completion of this coning cycle, the change ΔE_c in orbital coning energy can be written as

$$\begin{aligned} \Delta E_c &= (m/2) \cdot (D \cdot \omega_2)^2 \cdot (\alpha_0^2 - \alpha_1^2) \\ \Delta E_c &= m \cdot (D \cdot \omega_2)^2 \cdot [(\alpha_0 + \alpha_1)/2] \cdot (\alpha_0 - \alpha_1) \\ \Delta E_c &= m \cdot (D \cdot \omega_2)^2 \cdot \alpha \cdot \Delta \alpha \end{aligned}$$

where $(\alpha_0 + \alpha_1)/2 = \alpha$, the *average* coning angle over this cycle,
and $\alpha_0 - \alpha_1 = \Delta \alpha > 0$, the *reduction* in coning angle due to damping.

We now hypothesize that, at least in **hyper-stable flight** in which no nutation needs damping, and for **dynamically stable bullets**, the average loss in “forward motion” kinetic energy ΔE_α over any coning cycle due to flying with an *average* aerodynamic angle-of-attack α **causes** the *average*

“frictional damping” decrease in orbital coning energy ΔE_c during that same coning cycle. That is to say, we are tentatively assuming that a small fraction e ($e \approx 0.0025$) of the yaw-drag of the dynamically stable rifle bullet directly causes the damping of its coning angle α in *steady-state, minimum coning angle, “hyper-stable” flight*. Then, these two energy losses must be proportionally related to each other.

If this hypothesis is true, we can set $\Delta E_c = e \cdot \Delta E_\alpha$ over any particular coning cycle, where the constant fraction e is greater than **zero** for a **dynamically stable** projectile, but not greater than **1.0**, and so that

$$m \cdot (D \cdot \omega_2)^2 \cdot \alpha \cdot \Delta \alpha = q \cdot S \cdot V \cdot \alpha^2 \cdot e \cdot C_{D_\alpha} \cdot T_2$$

or, dividing through by $\alpha^2 > 0$ and by $[m \cdot (D \cdot \omega_2)^2]$,

$$(\Delta \alpha) / \alpha = (T_2) \cdot [q \cdot S \cdot V \cdot e \cdot C_{D_\alpha}] / [m \cdot (D \cdot \omega_2)^2].$$

We recognize this expression as having the form of the classic exponential damping of the coning angle α :

$$\alpha(t) = \alpha_0 \cdot \exp[-\lambda_2 \cdot t]$$

with $\lambda_2 = [q \cdot S \cdot V \cdot e \cdot C_{D_\alpha}] / [m \cdot (D \cdot \omega_2)^2]$.

If we replace the coning period T_2 with a small increment in time dt , and replace $\Delta \alpha$ per coning cycle with a small decrement $-d\alpha$ in α , then in the limit as dt approaches zero, this expression becomes

$$d\alpha / \alpha = -\lambda_2 \cdot dt$$

After integrating both sides from **0** to t ,

$$\ln[\alpha(t) / \alpha(0)] = -\lambda_2 \cdot t$$

Or, after exponentiating

$$\alpha(t) = \alpha_0 \cdot \exp[-\lambda_2 \cdot t] \quad [\text{QED}].$$

Thus, we have derived the long-accepted aeroballistic exponential damping relationship from the basic physics of our hypothesis that a portion e of the yaw-drag kinetic energy loss causes the damping of the coning angle for dynamically stable bullets in hyper-stable flight. [In “hyper-stable” flight in

flat-firing, the spin-stabilized projectile is flying with an aerodynamic angle-of-attack α equal to the change in flight path angle Φ caused only by gravity during each half coning cycle. This coning angle α is typically less than **0.1 degrees**, and produces negligible yaw-drag.]

If only a small fraction e ($0 < e \leq 1.0$) of this yaw-drag induced kinetic energy loss is actually responsible for frictional damping of the coning angle $\alpha(t)$ for a dynamically stable bullet, we accommodate that simply by using $e \cdot CD_\alpha$ in the above expression for λ_2 :

$$\lambda_2 = [q \cdot S \cdot V \cdot e \cdot CD_\alpha] / [m \cdot (D \cdot \omega_2)^2].$$

The **initial coning rate** ω_2 has been shown¹ always to be **directly proportional** to the selected rifling twist rate n given in calibers per turn, and none of the other parameters in this expression depend upon the selected rifling twist-rate n . [While n is used in evaluating e from data measurements for any given bullet, e is subsequently considered as a bullet constant.]

Therefore, the slow-mode damping factor λ_2 as formulated above varies **inversely with the square of n** . Smaller values of n produce larger damping factors λ_2 and quicker “damping out” of any initial coning angle α_0 and, thus, produce significantly less total kinetic energy loss over a long flight due to any given **initial yaw attitude** α_0 .

It is worth pointing out here that if $\lambda_2 = f_2$, the coning rate in hertz, the coning angle $\alpha(t)$ will be **critically damped** and reduced by a factor of **0.368** during each early coning cycle.

The total forward velocity ΔV lost by the rifle bullet due to yaw-drag while any initial coning angle α_0 is being damped out to **insignificance** can be well approximated from

$$\Delta KE = (m/2) \cdot (\Delta V)^2$$

$$\Delta KE = q \cdot S \cdot CD_\alpha \cdot V_0 \cdot \alpha_0^2 \int \exp[-2 \cdot \lambda_2 \cdot t] dt$$

$$\Delta KE = q \cdot S \cdot CD_\alpha \cdot V_0 \cdot \alpha_0^2 / (2 \cdot \lambda_2).$$

Here, we are integrating the drag-force component due to yaw in linear aeroballistics from $t = 0$ at the beginning of ballistic flight to some arbitrary later time T when minimum coning angle “hyper-stable” flight has been

achieved. For a dynamically stable rifle bullet fired from a very quick-twist ($n \approx 20$) barrel, T is just a very few early coning periods T_2 , and use of **initial values** is justified in this approximation.

From this expression, the **total retardation** of the bullet over a long-distance flight ΔV due to early yaw-drag is **directly proportional to the first power of n** , the rifling twist-rate in calibers per turn. This yaw-drag velocity loss occurs very early in ballistic flight and is never regained in flat firing; so, it *reduces* the **Ballistic Coefficient (BC)** measured for each bullet depending upon its particular initial non-zero yaw disturbance α_0 , which is consistent with our test-firing results.

When evaluating the coefficient of drag for exactly nose-forward flight CD_0 for any spin-stabilized projectile from measurements of a series of shots, only those shots having the **smallest** times-of-flight to the target, after correcting for any variation in launch speeds, are the most truly representative. Variable yaw-drag, caused by even small random yaw destabilization during launch, can only increase measured times-of-flight.

Explanatory Note 1:

The initial spin-rate f of the fired bullet in hertz is equal to the muzzle speed V_0 of the bullet in feet per second divided by the twist-rate ($n \cdot d$ in feet) of the rifling in the firing barrel.

From the Tri-Cyclic Theory, the initial coning rate ω_2 of the bullet in radians per second is then given by

$$\omega_2 = 2\pi \cdot (I_x/I_y) \cdot [V_0/(n \cdot d)] / (R + 1)$$

where

I_x, I_y = Second moments of inertia of the bullet about its crossed principal axes

n = Twist-rate in calibers per turn

d = Caliber in feet

$R = \omega_1/\omega_2 = f_1/f_2$ = Gyroscopic stability ratio, such that

$Sg = (R + 1)^2 / (4 \cdot R)$ = Gyroscopic stability.

According to information from a 1989 paper by Beat P. Kneubuehl of Thun, Switzerland, the initial value of $(R + 1)$ can be estimated for a simple cone-on-cylinder projectile model as

$$R + 1 = (l_x/l_y) * \{ [32\pi * l_x] / [n^2 * \rho * d^5 * CM_\alpha] \}$$

where

ρ = Air density in slugs per cubic foot

CM_α = Overturning moment coefficient in aeroballistics.

Both Sg and R vary *inversely with the square of twist-rate n* .

This is consistent with the classic aeroballistic definition of $Sg = P^2/(4*M)$ in terms of the auxiliary parameters P and M . P is a generalized spin-rate parameter (varying with n^{-1}) and M is a generalized overturning moment parameter as they are used in aeroballistics.

So, the initial coning rate ω_2 is *directly proportional to the selected twist-rate n* for the rifling.

For example, increasing n from **20** to **24 calibers per turn** (i.e., slowing the rifling twist-rate by **20-percent**) would increase the initial coning rate ω_2 of the rifle bullet by **20-percent**.

Importantly, this slowing of the twist-rate of the rifling by **20-percent** would simultaneously

(1) **increase** the total retardations ΔV caused by any random initial yaw disturbances α_0 by **20-percent** by decreasing the value of the slow-mode damping factor λ_2 by **44-percent**, and

(2) **decrease** the shot-group spread on the target by **20-percent** caused by aerodynamic jump deflections of the trajectory due to these same randomly oriented initial yaw disturbances, by reducing by **20-percent** the integration time $T_2/2$ over which the cross-track impulse is accumulated.